# Integral Equation Methods for Numerical Solutions of Partial Differential Equations

#### Zewen Shen

University of Toronto

August, 2020

This work is licensed under the Creative Commons Attribution-NonCommercial-ShareAlike 2.5 Canada License.



By the end of this talk, participants should be able to:

- Be more familiar with the development of numerical methods for PDEs.
- Output of the workflow of integral equation methods.
- Know the advantages and disadvantages of integral equation methods, and use it appropriately in future research.
- Show a subset of active problems arising from this field.

Accuracy

- Accuracy
- Speed

- Accuracy
- Speed
- Memory

- Accuracy
- Speed
- Memory

These requirements don't hold universally.

• Finding a good integral equation formulation requires work.

- Finding a good integral equation formulation requires work.
- There exist different integral equation formulations for a same PDE.

- Finding a good integral equation formulation requires work.
- There exist different integral equation formulations for a same PDE. What's the advantage of the formulations?

#### Definition

Laplace's equation with Dirichlet boundary condition

$$-\nabla^2 u(x) = 0 \quad \text{for } x \in \Omega \tag{1}$$
$$u(x) = f(x) \quad \text{for } x \in \partial\Omega \tag{2}$$

We want to find a solution to it that takes the form

$$u(x) = \int_{\partial\Omega} d(x,y)\sigma(y)ds(y), \quad x \in \Omega,$$

where d(x, y) is the normal derivative to the Green's function of Laplace's equation at the boundary point y $(d(x, y) = n(y) \cdot \nabla_y \phi(x - y) = \frac{n(y) \cdot (x - y)}{2\pi |x - y|^2}$  and n(y) is the unit length inwards pointing normal to  $\partial\Omega$  at  $y \in \partial\Omega$ ).

We want to find a solution to it that takes the form

$$u(x) = \int_{\partial\Omega} d(x,y)\sigma(y)ds(y), \quad x \in \Omega$$



$$u(x) = \int_{\partial\Omega} d(x,y)\sigma(y)ds(y), \quad x \in \Omega$$

### Questions

• Q: Why can we assume that the solution is written in this weird integral form?

$$u(x) = \int_{\partial\Omega} d(x,y)\sigma(y)ds(y), \quad x \in \Omega$$

### Questions

- Q: Why can we assume that the solution is written in this weird integral form?
- Q: Why do we pick this specific kernel function d(x, y) instead of others?

$$u(x) = \int_{\partial\Omega} d(x,y)\sigma(y)ds(y), \quad x \in \Omega$$

#### Summary

• Compute the correct  $\sigma$  such that the integral above matches the boundary condition.

$$u(x) = \int_{\partial\Omega} d(x,y)\sigma(y)ds(y), \quad x \in \Omega$$

#### Summary

- Compute the correct  $\sigma$  such that the integral above matches the boundary condition.
- Once σ is known, evaluate the PDE solution u(x) everywhere inside the domain by computing the line integral.

$$u(x) = \int_{\partial\Omega} d(x,y)\sigma(y)ds(y), \quad x \in \Omega$$

#### Summary

- Compute the correct  $\sigma$  such that the integral above matches the boundary condition.
- Once σ is known, evaluate the PDE solution u(x) everywhere inside the domain by computing the line integral.

How to solve the density function  $\sigma(y)$  defined on the boundary?

Apply the boundary condition: u(x) = f(x) for  $x \in \partial \Omega$ ,

How to solve the density function  $\sigma(y)$  defined on the boundary?

Apply the boundary condition: u(x) = f(x) for  $x \in \partial\Omega$ , or more rigorously:  $\lim_{x' \to x, x' \in \Omega} u(x') = f(x)$  for  $x \in \partial\Omega$ 

How to solve the density function  $\sigma(y)$  defined on the boundary?

Apply the boundary condition: u(x) = f(x) for  $x \in \partial\Omega$ , or more rigorously:  $\lim_{x' \to x, x' \in \Omega} u(x') = f(x)$  for  $x \in \partial\Omega$ 

#### Remarks

Our kernel function d(x, y) has a singularity as  $x \to y$ , which leads to an extra term as we take the limit

$$\lim_{x'\to x,x'\in\Omega}u(x')=-\frac{1}{2}\sigma(x)+\int_{\partial\Omega}d(x,y)\sigma(y)ds(y).$$

How to solve the density function  $\sigma(y)$  defined on the boundary?

Apply the boundary condition: u(x) = f(x) for  $x \in \partial\Omega$ , or more rigorously:  $\lim_{x' \to x, x' \in \Omega} u(x') = f(x)$  for  $x \in \partial\Omega$ 

#### Remarks

Our kernel function d(x, y) has a singularity as  $x \rightarrow y$ , which leads to an extra term as we take the limit

$$\lim_{x'\to x,x'\in\Omega}u(x')=-\frac{1}{2}\sigma(x)+\int_{\partial\Omega}d(x,y)\sigma(y)ds(y).$$

Therefore, we got the integral equation that  $\sigma(x)$  satisfies:

$$-rac{1}{2}\sigma(x)+\int_{\partial\Omega}d(x,y)\sigma(y)ds(y)=f(x),\quad x\in\partial\Omega.$$

#### References

- J. Levandosky, Notes on potential theory, Stanford University, 2003.
- P.G. Martinsson, Direct Solvers for Integral Equations, CBMS/NSF Conference on Fast Direct Solvers, 2014.
- P.G. Martinsson, Fast Direct Solvers for Elliptic PDEs, SIAM, 2019

## Discretization of the integral equation

The density function  $\sigma(x)$  satisfies a Fredholm integral equation of second kind:

$$-rac{1}{2}\sigma(x)+\int_{\partial\Omega}d(x,y)\sigma(y)ds(y)=f(x),\quad x\in\partial\Omega.$$

We discretize the integral equation by replacing the integral with a representative weighted sum, i.e., a selected quadrature rule:

$$\int_{\partial\Omega} \phi(y) ds(y) = \sum_{j=1}^{N} w_j \phi(x_j), \quad \text{for a certain class of } \phi.$$

## Discretization of the integral equation

The density function  $\sigma(x)$  satisfies a Fredholm integral equation of second kind:

$$-rac{1}{2}\sigma(x)+\int_{\partial\Omega}d(x,y)\sigma(y)ds(y)=f(x),\quad x\in\partial\Omega.$$

We discretize the integral equation by replacing the integral with a representative weighted sum, i.e., a selected quadrature rule:

$$\int_{\partial\Omega} \phi(y) ds(y) = \sum_{j=1}^{N} w_j \phi(x_j), \quad \text{for a certain class of } \phi.$$

Then the discretized integral equation becomes a  $N \times N$  linear system

$$-\frac{1}{2}\sigma(x_i)+\sum_{j=1}^N w_j d(x_i,x_j)\sigma(x_j)=f(x_i), \quad i=1,2,\ldots,N.$$

#### Examples of common quadrature rules

- Trapezoidal rule: Converge super-algebraically when the boundary condition f is  $C^{\infty}$  and periodic, plus the boundary of the domain is smooth. Easy to code.
- Composite Gaussian quadrature: Customized convergence rate depending on the selected order of quadrature. Allows for local refinement, which is good for domains with corners and edges.
- Singular integral quadrature: Necessary when we need to evaluate the integral of a singular kernel. We don't need it in our example, but it's required for more sophisticated problems, e.g., Helmholtz equation.

## Discretization of the integral equation

To write the discretized integral equation

$$-\frac{1}{2}\sigma(x_i) + \sum_{j=1}^N w_j d(x_i, x_j)\sigma(x_j) = f(x_i)$$

as a N imes N linear system, we define vectors  $\sigma, f \in R^n$  via

$$f(i) = \sqrt{w_i}f(x_i), \quad \sigma(i) = \sqrt{w_i}\sigma(x_i),$$

and define matrix  $A \in R^{N \times N}$  via

$$A(i,j) = \sqrt{w_i} d(x_i, x_j) \sqrt{w_j} - \frac{1}{2} \delta_{i,j}.$$
(3)

Finally, we get a linear system  $A\sigma = f$  to solve.

### Why do we need the $\sqrt{w}$ weighting scheme?

Such weighting scheme allows the matrix to capture the  $L^2$  behaviour of an integral operator rather than its pointwise behavior, which leads to a better-conditioned linear system due to some intrinsic property of our integral equation.

Zewen Shen (UToronto)

#### References

- J Bremer, Z Gimbutas, V Rokhlin, A nonlinear optimization procedure for generalized Gaussian quadratures, SIAM Journal on Scientific Computing, 2010
- S Hao, AH Barnett, PG Martinsson, P Young, High-order accurate methods for Nyström discretization of integral equations on smooth curves in the plane, Advances in Computational Mathematics, 2014
- J. Bremer, A fast direct solver for the integral equations of scattering theory on planar curves with corners, Journal of Computational Physics, 2012
- J. Bremer, A Nyström method for weakly singular integral operators on surfaces, Journal of Computational Physics, 2012

Let's take a close look at the matrix A where

$$A(i,j) = \sqrt{w_i} d(x_i, x_j) \sqrt{w_j} - \frac{1}{2} \delta_{i,j}.$$
(4)

Unlike the usual sparse matrices arising from the common numerical PDE methods, our matrix is dense and rank-structured! This tells us...

Let's take a close look at the matrix A where

$$A(i,j) = \sqrt{w_i} d(x_i, x_j) \sqrt{w_j} - \frac{1}{2} \delta_{i,j}.$$
(4)

Unlike the usual sparse matrices arising from the common numerical PDE methods, our matrix is dense and rank-structured! This tells us...

• Conventional matrix solvers will be slow.

Let's take a close look at the matrix A where

$$A(i,j) = \sqrt{w_i} d(x_i, x_j) \sqrt{w_j} - \frac{1}{2} \delta_{i,j}.$$
(4)

Unlike the usual sparse matrices arising from the common numerical PDE methods, our matrix is dense and rank-structured! This tells us...

- Conventional matrix solvers will be slow.
- The matrix requires more storage as its columns can't be stored sparsely.

$$A(i,j) = \sqrt{w_i}d(x_i,x_j)\sqrt{w_j} - rac{1}{2}\delta_{i,j}$$
 ?

$$A(i,j) = \sqrt{w_i}d(x_i,x_j)\sqrt{w_j} - \frac{1}{2}\delta_{i,j}$$
 ?

• Each entry represents an interaction between point  $x_i$  and  $x_j$ .

$$A(i,j) = \sqrt{w_i} d(x_i, x_j) \sqrt{w_j} - \frac{1}{2} \delta_{i,j} \quad ?$$

- Each entry represents an interaction between point  $x_i$  and  $x_j$ .
- Such interactions decay in magnitude over distances...

$$A(i,j) = \sqrt{w_i} d(x_i, x_j) \sqrt{w_j} - \frac{1}{2} \delta_{i,j} \quad ?$$

- Each entry represents an interaction between point  $x_i$  and  $x_j$ .
- Such interactions decay in magnitude over distances...
- Maybe we can set the interaction to 0 if two points are far from each other enough?

$$A(i,j) = \sqrt{w_i} d(x_i, x_j) \sqrt{w_j} - \frac{1}{2} \delta_{i,j} \quad ?$$

- Each entry represents an interaction between point  $x_i$  and  $x_j$ .
- Such interactions decay in magnitude over distances...
- Maybe we can set the interaction to 0 if two points are far from each other enough?

However, while the magnitudes of these interactions decay only slowly, the amount of information that they carry decays enormously fast. We call such interaction low-rank.

## Rank-structured matrix

While the magnitudes of these fields decay only slowly, the amount of information that they carry decays enormously fast. We call such interaction low-rank.

Definition: Given an error tolerance  $\epsilon$ , if A is a  $m \times n$  low-rank matrix with numerical rank k, then we have  $U \in R^{m \times k}$  and  $V \in R^{k \times n}$ 

 $A \approx UV$ ,

and the numerical rank k is the smallest integer such that  $|A - UV|_2 < \epsilon$ : numrank $(A, \epsilon) = \min\{k : \exists U \in R^{m \times k}, V \in R^{k \times n}, s.t., |A - UV|_2 < \epsilon\}.$ 

# Rank-structured matrix

While the magnitudes of these fields decay only slowly, the amount of information that they carry decays enormously fast. We call such interaction low-rank.

Definition: Given an error tolerance  $\epsilon$ , if A is a  $m \times n$  low-rank matrix with numerical rank k, then we have  $U \in R^{m \times k}$  and  $V \in R^{k \times n}$ 

### $A \approx UV$ ,

and the numerical rank k is the smallest integer such that  $|A - UV|_2 < \epsilon$ : numrank $(A, \epsilon) = \min\{k : \exists U \in R^{m \times k}, V \in R^{k \times n}, s.t., |A - UV|_2 < \epsilon\}.$ 

## Remarks

The numerical rank of a low rank matrix is usually far lower than the actual rank, i.e.,  $k \ll \min(m, n)$ . What happens to the storage?

Zewen Shen (UToronto)

# Rank-structured matrix

While the magnitudes of these fields decay only slowly, the amount of information that they carry decays enormously fast. We call such interaction low-rank.

Definition: Given an error tolerance  $\epsilon$ , if A is a  $m \times n$  low-rank matrix with numerical rank k, then we have  $U \in R^{m \times k}$  and  $V \in R^{k \times n}$ 

### $A \approx UV$ ,

and the numerical rank k is the smallest integer such that  $|A - UV|_2 < \epsilon$ : numrank $(A, \epsilon) = \min\{k : \exists U \in R^{m \times k}, V \in R^{k \times n}, s.t., |A - UV|_2 < \epsilon\}.$ 

## Remarks

The numerical rank of a low rank matrix is usually far lower than the actual rank, i.e.,  $k \ll \min(m, n)$ . What happens to the storage? What happens to the matvec operation?

Zewen Shen (UToronto)

Integral Equation Methods

# Rank-structured matrix

While the magnitudes of these fields decay only slowly, the amount of information that they carry decays enormously fast. We call such interaction low-rank.

Definition: Given an error tolerance  $\epsilon$ , if A is a  $m \times n$  low-rank matrix with numerical rank k, then we have  $U \in R^{m \times k}$  and  $V \in R^{k \times n}$ 

### $A \approx UV$ ,

and the numerical rank k is the smallest integer such that  $|A - UV|_2 < \epsilon$ : numrank $(A, \epsilon) = \min\{k : \exists U \in R^{m \times k}, V \in R^{k \times n}, s.t., |A - UV|_2 < \epsilon\}.$ 

## Remarks

The numerical rank of a low rank matrix is usually far lower than the actual rank, i.e.,  $k \ll \min(m, n)$ . What happens to the storage? What happens to the matvec operation?

Zewen Shen (UToronto)

Integral Equation Methods

How do we compute a low-rank approximation of the matrix?

## Examples

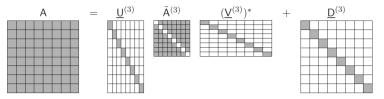
- SVD: Optimal! (by Eckart-Young-Mirsky Theorem) But...
- Randomized range finder: Fast. Receiving more and more attentions nowadays. However, its advantage is not obvious compared to the next example in our application...
- Interpolatory decomposition: Widely used by the integral equation method researchers! Nearly optimal, reasonably fast, no need to construct the whole matrix A when integrated into our algorithm (by exploiting the analytical property of our kernel...)

# References

- A. Kloeckner, Tools for Low-Rank Linear Algebra, Fast Algorithms and Integral Equation Methods, Fall 2019
- P.G. Martinsson, Fast Direct Solvers for Elliptic PDEs, SIAM, 2019

# Low-rank approximation of our matrix

Low-rank approximation of the first level (cited from [Gillman, et al])



Final Low-rank approximation (cited from [Gillman, Young, Martinsson])

$$(4.4) \qquad A = \underline{U}^{(3)} (\underline{U}^{(2)} (\underline{U}^{(1)} \underline{B}^{(0)} (\underline{V}^{(1)})^* + \underline{B}^{(1)}) (\underline{V}^{(2)})^* + \underline{B}^{(2)}) (\underline{V}^{(3)})^* + \underline{D}^{(3)}.$$
  
The block structure of the right hand side of (4.4) is:  
$$\underline{\underline{U}^{(3)}} \underline{\underline{U}^{(2)}} \underline{\underline{U}^{(1)}} \underline{\underline{B}^{(0)}} (\underline{\underline{V}^{(1)}})^* \underline{\underline{B}^{(1)}} (\underline{\underline{V}^{(2)}})^* \underline{\underline{B}^{(2)}} (\underline{\underline{V}^{(3)}})^* \underline{\underline{D}^{(3)}}.$$

• Time complexity of the factorization:  $O(Nk^2)$ . No need for first constructing the whole matrix A.

$$(4.4) \qquad A = \underline{U}^{(3)} (\underline{U}^{(2)} (\underline{U}^{(1)} \underline{B}^{(0)} (\underline{V}^{(1)})^* + \underline{B}^{(1)}) (\underline{V}^{(2)})^* + \underline{B}^{(2)}) (\underline{V}^{(3)})^* + \underline{D}^{(3)}.$$
  
The block structure of the right hand side of (4.4) is:  
$$\underline{\underline{U}^{(3)}} \underline{U}^{(2)} \underline{U}^{(1)} \underline{B}^{(0)} (\underline{V}^{(1)})^* \underline{B}^{(1)} (\underline{V}^{(2)})^* \underline{B}^{(2)} (\underline{V}^{(3)})^* \underline{D}^{(3)}.$$

- Time complexity of the factorization:  $O(Nk^2)$ . No need for first constructing the whole matrix A.
- Time complexity of matvec: O(kN).

$$A(i,j) = \sqrt{w_i}d(x_i, x_j)\sqrt{w_j} - \frac{1}{2}\delta_{i,j}$$

$$A(i,j) = \sqrt{w_i} d(x_i, x_j) \sqrt{w_j} - \frac{1}{2} \delta_{i,j}$$

• Analytical expansion vs matrix factorization.

$$A(i,j) = \sqrt{w_i}d(x_i, x_j)\sqrt{w_j} - \frac{1}{2}\delta_{i,j}$$

- Analytical expansion vs matrix factorization.
- Strong admissibility vs weak admissibility.

### References

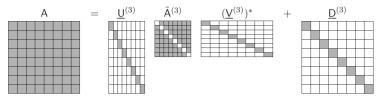
- A. Gillman, P. Young, P.G. Martinsson, A direct solver with O(N) complexity for integral equations on one-dimensional domains, Front. Math. China, 7 (2012), pp. 217–247.
- R. Beatson, L. Greengard, A short course on fast multipole methods
- L. Greengard, V. Rokhlin, A fast algorithm for particle simulations, Journal of Computational Physics, 1987

• Multiple RHS.

- Multiple RHS.
- Independent of condition number.

- Multiple RHS.
- Independent of condition number.

Next question: How to compute the inverse of our factorized matrix?



There's a name for such factorization...

$$(4.4) \qquad A = \underline{U}^{(3)} \left( \underline{U}^{(2)} \left( \underline{U}^{(1)} \underline{B}^{(0)} \left( \underline{V}^{(1)} \right)^* + \underline{B}^{(1)} \right) \left( \underline{V}^{(2)} \right)^* + \underline{B}^{(2)} \right) \left( \underline{V}^{(3)} \right)^* + \underline{D}^{(3)}.$$
  
The block structure of the right hand side of (4.4) is:  
$$\underline{U}^{(3)} \quad \underline{U}^{(2)} \quad \underline{U}^{(1)} \quad \underline{B}^{(0)} \quad (\underline{V}^{(1)})^* \quad \underline{B}^{(1)} \quad (\underline{V}^{(2)})^* \quad \underline{B}^{(2)} \quad (\underline{V}^{(3)})^* \quad \underline{D}^{(3)}.$$

• Sherman-Morrison-Woodbury formula

$$(4.4) \qquad A = \underline{U}^{(3)} (\underline{U}^{(2)} (\underline{U}^{(1)} \underline{B}^{(0)} (\underline{V}^{(1)})^* + \underline{B}^{(1)}) (\underline{V}^{(2)})^* + \underline{B}^{(2)}) (\underline{V}^{(3)})^* + \underline{D}^{(3)}.$$
  
The block structure of the right hand side of (4.4) is:  
$$\underline{U}^{(3)} \underline{U}^{(2)} \underline{U}^{(1)} \underline{B}^{(0)} (\underline{V}^{(1)})^* \underline{B}^{(1)} (\underline{V}^{(2)})^* \underline{B}^{(2)} (\underline{V}^{(3)})^* \underline{D}^{(3)}.$$

- Sherman-Morrison-Woodbury formula
- Time complexity of computing inverse:  $O(Nk^2)$

$$(4.4) \qquad A = \underline{U}^{(3)} (\underline{U}^{(2)} (\underline{U}^{(1)} \underline{B}^{(0)} (\underline{V}^{(1)})^* + \underline{B}^{(1)}) (\underline{V}^{(2)})^* + \underline{B}^{(2)}) (\underline{V}^{(3)})^* + \underline{D}^{(3)}.$$
  
The block structure of the right hand side of (4.4) is:  
$$\underline{U}^{(3)} \underline{U}^{(2)} \underline{U}^{(1)} \underline{B}^{(0)} (\underline{V}^{(1)})^* \underline{B}^{(1)} (\underline{V}^{(2)})^* \underline{B}^{(2)} (\underline{V}^{(3)})^* \underline{D}^{(3)}.$$

- Sherman-Morrison-Woodbury formula
- Time complexity of computing inverse:  $O(Nk^2)$
- Time complexity of applying the inverse: O(Nk)

## References

- J. Bremer, A. Gillman, P.G. Martinsson, A high-order accurate accelerated direct solver for acoustic scattering from surfaces, BIT Numerical Mathematics, 2015
- A. Gillman, P. Young, P.G. Martinsson, A direct solver with O(N) complexity for integral equations on one-dimensional domains, Front. Math. China, 7 (2012), pp. 217–247.
- P.G. Martinsson, Fast Direct Solvers for Elliptic PDEs, SIAM, 2019
- Z. Shen, lap2d: A Fast Direct Dense Solver for 2-D Laplace's Equation, 2020

• High-order convergence rate plus reduction of dimensionality.

- High-order convergence rate plus reduction of dimensionality.
- Be able to achieve machine precision accuracy.

- High-order convergence rate plus reduction of dimensionality.
- Be able to achieve machine precision accuracy.
- Works well for exterior problems (acoustics, electromagnetic scattering).

- High-order convergence rate plus reduction of dimensionality.
- Be able to achieve machine precision accuracy.
- Works well for exterior problems (acoustics, electromagnetic scattering).
- Capability of exploiting all kinds of analytical properties of PDEs. Nice combination between analysis and numerical linear algebra!

Disadvantages

• Harder to code.

# Disadvantages

- Harder to code.
- Special considerations need to be done when the problem has body load, singular kernels, a domain with corners, mixed boundary condition, etc.

# Questions?

Thanks to prof. Kirill Serkh and Xinyue Lu for providing suggestions to the talk.